

A Modified Boundary Element Method for Seismic Modeling

> PhD Student: Yi Tao Advisor: Mrinal Sen February, 2010



SCHOOL OF GEOSCIENCES

Outline

- Introduction
- Modified Boundary Element Method
- Benchmark test
- Wave scattering in the frequency domain
- Synthetic seismogram examples
- Conclusions

 Three major numerical simulation methods can be used for seismic modeling: FD (Finite Difference), FEM (Finite Element Method) and BEM (Boundary Element Method)

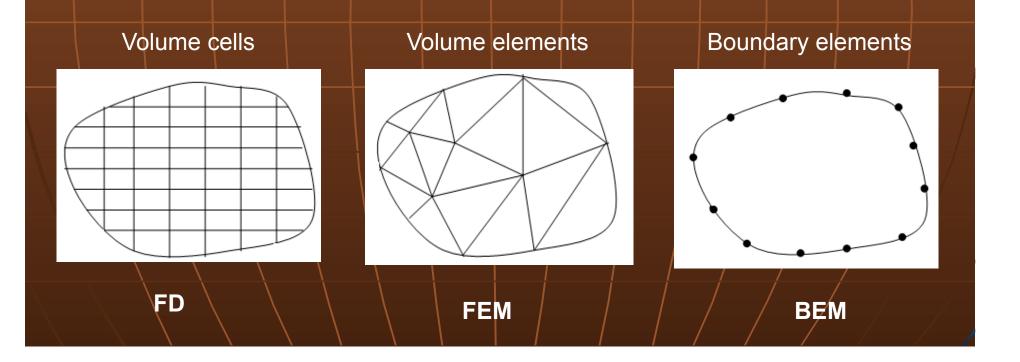
•FD is the most widely used method in seismic modeling and migration.

 Strong impedance change among irregular interfaces, surface topographies and volume heterogeneities challenge the accuracy of FD.

 FEM is another major method, compared to FD, it has the advantage of accuracy for the above geological structures.

• BEM (Boundary Element Method) (Sanchez-Sesma, 1984, Bouchon, 1995, Fu, 2002) has the same advantage of accuracy of FEM.

 Compared to FEM, BEM only requires grids on the boundary, has the advantage of less data preparation, but it requires that the media inside the domain is homogeneous.



Comparison of the three methods (Ordinary FD, Ordinary FEM, Direct BEM)

	FD	FEM	BEM
Domain	Whole volume	Whole volume	Boundary
Mesh	Regular mesh	Irregular mesh	Irregular mesh
Boundary conditions	Needs special treatment on	Free surface conditions	Free surface and inner interface
	curved structures	automatically satisfied.	conditions automatically satisfied.
Cost	Fast	Expensive (sparse global matrix)	Expensive (usually full global matrix)
Heterogeneity	ОК	OK	Problematic

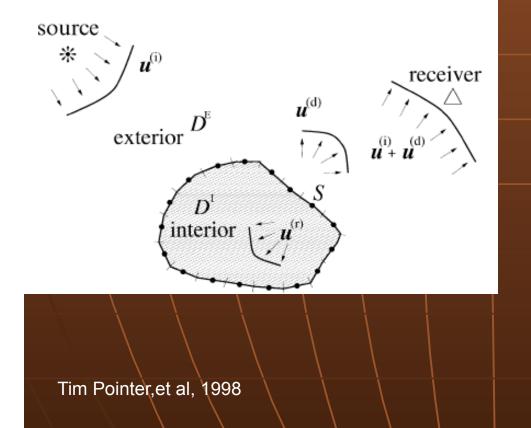
• We use a modified BEM method which can simulate volume heterogeneities.

•Compared to BEM, it adds a volume scattering term. So it has all the advantages of BEM, only requires an additional storage of the volume scattering data.

•BEM can only be applied to piecewise heterogeneous media, if we know how to describe the scattering term, the modified method can be applied to general heterogeneous media.

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 $u(\mathbf{r}) = u^{i}(\mathbf{r}) + u^{d}(\mathbf{r})$

The total wavefield recorded at a receiver which passes through a scatter can be written as the incident wavefield plus a scattered wavefield

Acoustic wave Equation in the Frequency Domain:

 $\nabla^2 u(\mathbf{r}) + k(\mathbf{r})^2 u(\mathbf{r}) = -s(\mathbf{r},\omega)$

Define a relative slowness perturbation

$$o(\mathbf{r}') = \rho(\mathbf{r})\mu_0(\mathbf{r})/\rho_0(\mathbf{r})\mu(\mathbf{r})-1$$

$$\nabla^2 u(\mathbf{r}) + k_0(\mathbf{r})^2 u(\mathbf{r}) = -s(\mathbf{r},\omega) - k_0(\mathbf{r})^2 o(\mathbf{r})u(\mathbf{r})$$

Generalized Lippmann-Schwinger Equation

$$\int_{\Gamma} \left[G(\mathbf{r},\mathbf{r}') \frac{\partial u(\mathbf{r}')}{\partial n} - u(\mathbf{r}') \frac{\partial G(\mathbf{r},\mathbf{r}')}{\partial n} \right] d\mathbf{r} + k_0^2 \int_{\Omega} o(\mathbf{r}') u(\mathbf{r}') G(\mathbf{r},\mathbf{r}') d\mathbf{r}' + S(\omega) G(\mathbf{r},\mathbf{r}_0) d\mathbf{r}' + S(\omega) G(\mathbf{r},\mathbf{r}_0) d\mathbf{r}' d\mathbf{r}'$$

$$=\begin{cases} u(\mathbf{r}) & \mathbf{r} \in \Omega \\ C(\mathbf{r})u(\mathbf{r}) & \mathbf{r} \in \Gamma \\ 0 & \mathbf{r} \notin \Omega \end{cases}$$

Boundary conditions for layered media:

1.Free surface, the displacement is zero

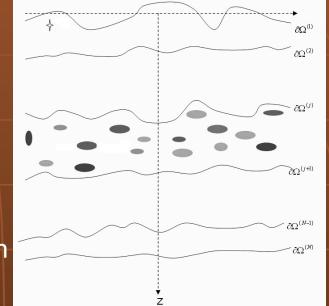
 $\frac{\partial u(\mathbf{r})}{\partial n} = 0 \qquad \mathbf{r} \in \partial \Omega^{\text{flock}}$

2.Inner interfaces: the displacement and traction are continuous along the interface

$$\mu^{(j)} \frac{\partial u_{-}^{(j)}(\mathbf{r})}{\partial n} = \mu^{(j+1)} \frac{\partial u_{+}^{j}(\mathbf{r})}{\partial n} \quad \mathbf{r} \in \partial \Omega^{j \in \mathcal{D}}$$

3. Infinite boundary: Sommerfield far field condition

$$\begin{cases} \lim_{\mathbf{r}\to\infty} u(\mathbf{r}) = 0\\ \lim_{\mathbf{r}\to\infty} \frac{\partial u(\mathbf{r})}{\partial r} = iK_0 u(\mathbf{r}) \end{cases}$$



For layered media:

$$\int_{\partial \Omega^{(j)}} \left[\left(\tilde{\mathcal{C}}_{j,w}^{j;\tilde{\mathbf{r}}}(\tilde{\mathbf{r}}', \tilde{\omega}) \frac{\partial u^{(j)}(\mathbf{r}')}{\partial n} - u^{(j)}(\mathbf{r}') \frac{\partial G^{(j)}(\mathbf{r}, \mathbf{r}')}{\partial n} \right] d\mathbf{r}' \tilde{\mathfrak{t}} \int_{\partial \Omega^{(j)}} \left[G^{\ell j}(\mathbf{r}', \omega) \frac{\partial u^{(j;\tilde{\mathbf{r}})}(\mathbf{r}')}{\partial n} - u^{(j-1)}(\mathbf{r}') \frac{\partial G^{(j)}(\mathbf{r}, \mathbf{r}')}{\partial n} \right] d\mathbf{r}'$$

$$+ \left[K_{0}^{(j)} \right]^{2} \int_{\Omega^{(j)}_{s'}} \mathcal{O}^{j;\tilde{\mathbf{r}}}(\tilde{\mathbf{r}}') \tilde{\mathfrak{s}}^{\ell j;\tilde{\mathbf{r}}}(\mathbf{r}') G^{(j)}(\mathbf{r}, \mathbf{r}') d\mathbf{r}' = \begin{cases} u^{(j)}(\mathbf{r}); \ i \ i \ \mathbf{r} \in \Omega^{(j)} \\ C^{(j)}(\mathbf{r}) u^{(j)}(\mathbf{r}); \ \mathbf{r} \in \partial \Omega^{(j)} \\ C^{(j+1)}(\mathbf{r}) u^{(j+1)}(\mathbf{r}); \ \mathbf{r} \in \partial \Omega^{(j+1)} \\ 0; \ i \ i \ \mathbf{r} \notin \overline{\Omega} \end{cases}$$

$$= 0 \quad \text{The media within a layer is homogeneous}$$

We use Green function in the background media

2D acoustic media:

$$G(\mathbf{r},\mathbf{r}_0) = \frac{i}{4} H_0^{(1)}(k|\mathbf{r}-\mathbf{r}_0|)$$

3D acoustic media:

$$G(\mathbf{r},\mathbf{r}_0) = \frac{1}{4\pi |\mathbf{r} - \mathbf{r}_0|} \exp(ik|\mathbf{r} - \mathbf{r}_0|)$$

Elastic media:

3

$$\begin{aligned} T_{ki} &= \frac{\partial G_{ki}(\mathbf{r}, \mathbf{r}_{0})}{\partial n} = \mu A \left\{ \left\{ \begin{array}{l} \delta_{kk}^{\perp} \frac{\partial^{2} \sigma_{0}}{\partial n} \cdot n_{k} r_{k} & \frac{\lambda}{\mu} n_{r} r_{k} \right\} \frac{\partial U_{1}}{\partial r} \\ &- \left\{ \delta_{kk}^{\perp} \frac{\partial^{2} \sigma_{0}}{\partial n} \cdot 2k_{k} r_{i} & n_{r} r_{k} - 2r_{k} r_{i} \frac{\partial r}{\partial n} + a \frac{\lambda}{\mu} + n_{r} r_{k} \right\} \frac{U}{r} \\ &- \left[2r_{k} r_{i} \frac{\partial r}{\partial n} + \frac{\lambda}{\mu} + n_{r} r_{k} \right] \frac{U_{2}}{r} \right\} \end{aligned}$$

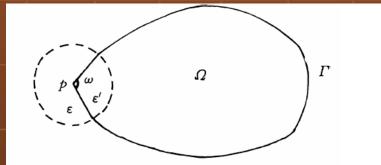
$$2D \text{ elastic:} \quad A = \frac{i}{4\mu} \quad U_{1} = H_{0}^{(1)}(k_{\beta}r) - \frac{1}{k_{\beta}r} H_{1}^{(1)}(k_{\beta}r) + \left(\frac{k_{\alpha}}{k_{\beta}}\right)^{2} \frac{1}{k_{\alpha}r} H_{1}^{(1)}(k_{\alpha}r) \\ &U_{2} = -H_{0}^{(1)}(k_{\beta}r) + \left(\frac{k_{\alpha}}{k_{\beta}}\right)^{2} \frac{1}{k_{\alpha}r} H_{1}^{(1)}(k_{\alpha}r) \\ &U_{2} = -H_{0}^{(1)}(k_{\beta}r) - \frac{1}{k_{\beta}r} - \left(\frac{1}{k_{\beta}r}\right)^{2} \left| \frac{e^{ik_{\beta}r}}{r} - \left(\frac{k_{\alpha}}{k_{\beta}}\right)^{2} \left| \frac{i}{k_{\alpha}r} - \left(\frac{1}{k_{\alpha}r}\right)^{2} \right| \frac{e^{ik_{\alpha}r}}{r} \\ &U_{2} = \left[1 + \frac{3i}{k_{\beta}r} - 3\left(\frac{1}{k_{\beta}r}\right)^{2} \right] \frac{e^{ik_{\beta}r}}{r} - \left(\frac{k_{\alpha}}{k_{\beta}}\right)^{2} \left[1 + \frac{3i}{k_{\alpha}r} - 3\left(\frac{1}{k_{\alpha}r}\right)^{2} \right] \frac{e^{ik_{\alpha}r}}{r} \end{aligned}$$

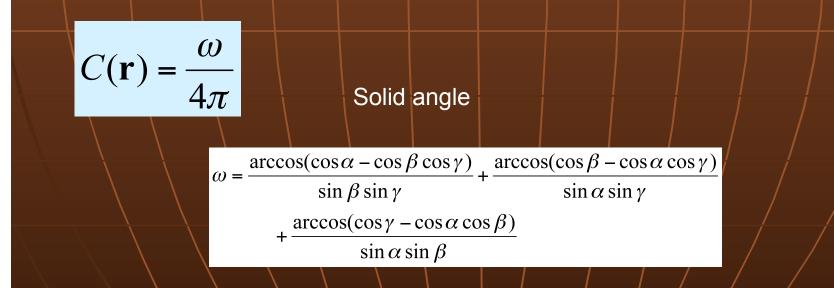
The coefficient C(r) in acoustic media: Smooth points: C(r)=0.5

2-D non-smooth points:

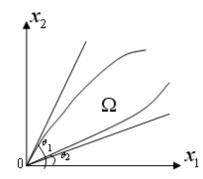
 $C(\mathbf{r}) = \frac{\omega}{2\pi}$

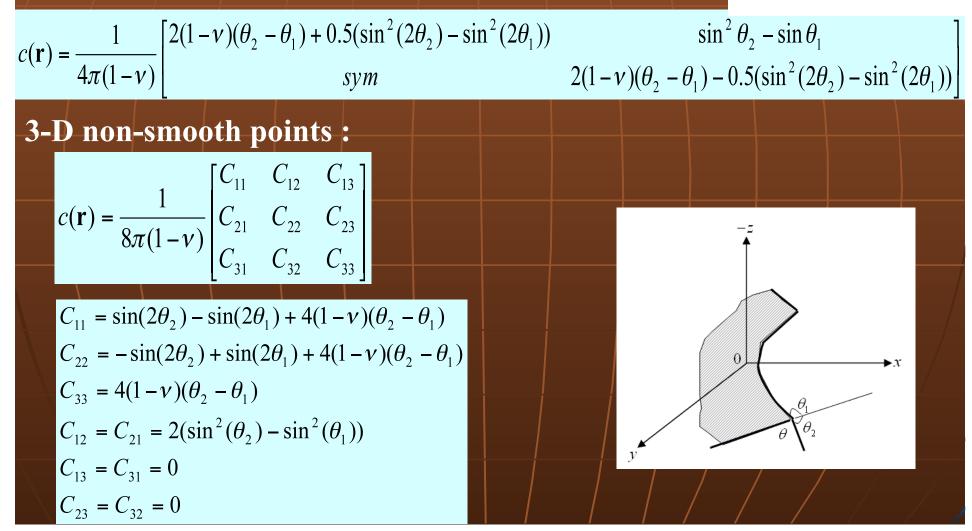
3-D non-smooth points :





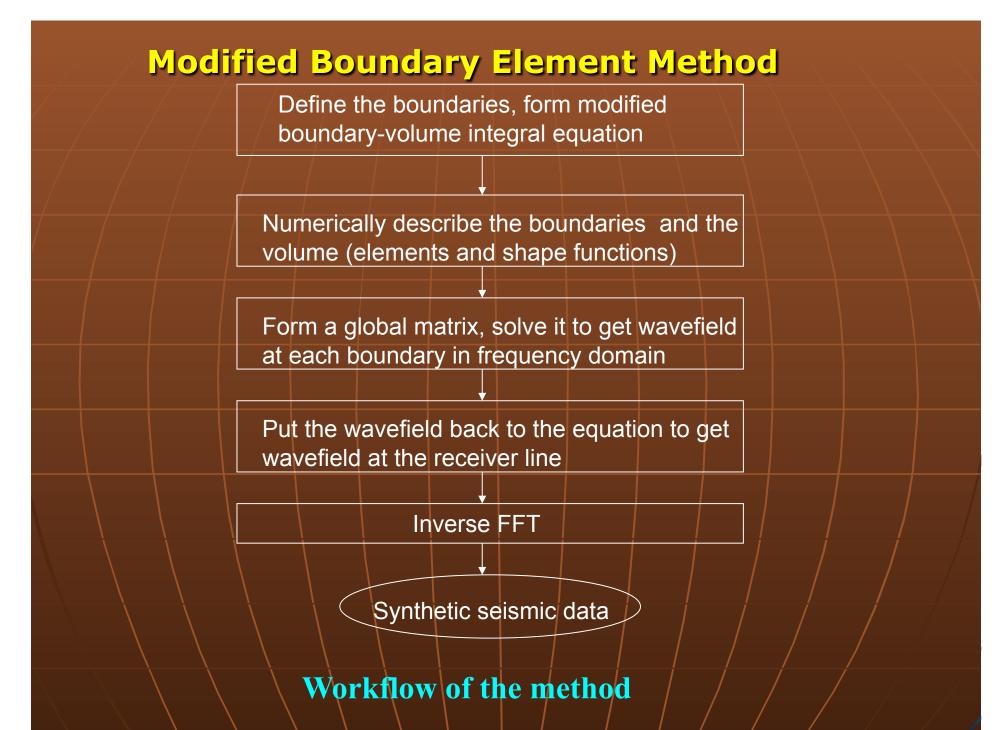
The coefficient C(r) in elastic media: Smooth points: C(r)=0.5 2-D non-smooth points:





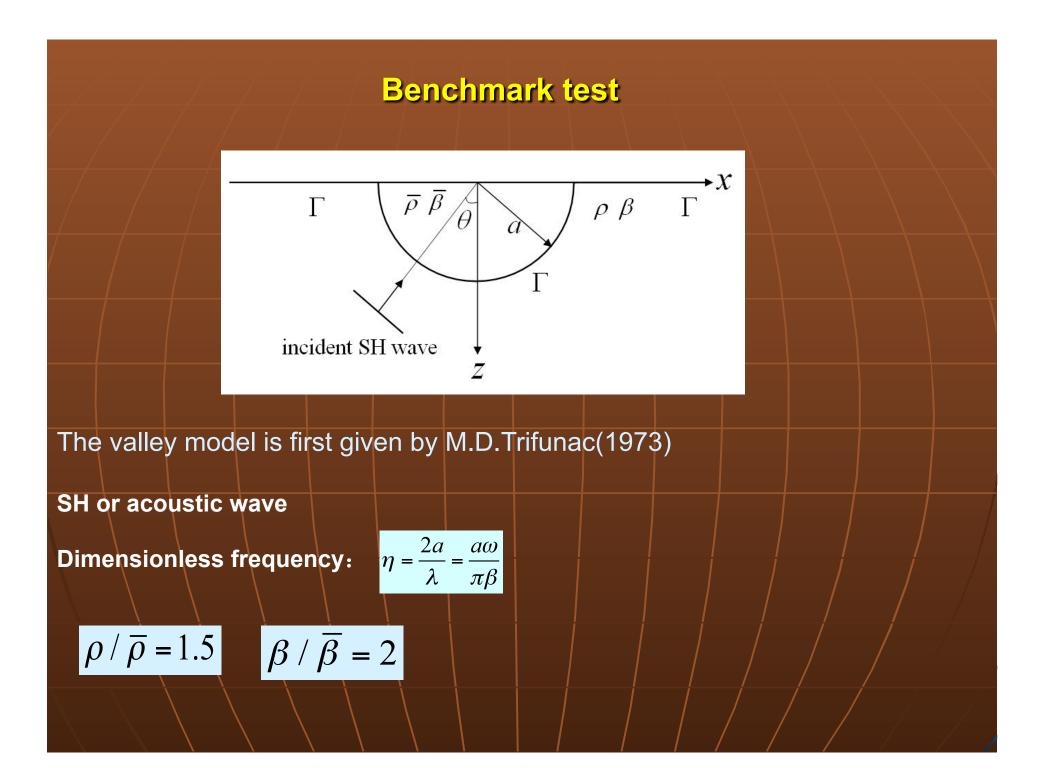
Modified Boundary Element Method Numerical implementation of the method **Boundary scattering (BEM):** $H^{i,\mathfrak{t}} \overset{\mathfrak{t}}{\mathbf{u}}^{(i)} \overset{\mathfrak{t}}{(\mathbf{r}_{i})} \overset{\mathfrak{t}}{-} \overset{\mathfrak{t}}{G}^{\sharp i,\mathfrak{t}} \overset{\mathfrak{t}}{\mathbf{v}}^{(i)} \overset{\mathfrak{t}}{(\mathbf{r}_{i})} \mathfrak{t} \overset{\mathfrak{t}}{H}^{i,2} \mathfrak{t}^{(i-1)}(\mathbf{r}_{i}) \ltimes G^{i,2} \mathfrak{t}^{(i-1)}(\mathbf{r}_{i})$ **Volume scattering :** 20 $K^{\mathfrak{t}} \mathbf{W}^{\mathfrak{t}} \mathbf{W}^{\mathfrak{t}}(\mathbf{r}_{i}^{\mathfrak{G}})$ $\mathbf{A}_{1}^{(i)}\mathbf{Q}^{\text{fife}} + \mathbf{A}_{2}^{(i)}\mathbf{Q}^{\text{fife}} + \mathbf{K}^{(i)}w^{(i)} = \boldsymbol{\delta}_{si}\mathbf{s}$ $\mathbf{A}_{1}^{(N \pounds \ 1)}$ « $m^{\text{f} NE+1}$ $\mathbf{A}_{2}^{(N)} \quad \mathbf{K}^{\mathbb{E}^{N}} \quad {}^{\mathbb{O}}\mathbf{A}_{1}^{(N)}$ $m^{(N)}$ 0 0 $\mathbf{A}_{2}^{(i+1)} \quad \mathbf{K}^{\mathfrak{t} \ i \ \mathfrak{l} \ \mathfrak{t}} \quad \mathbf{A}_{1}^{(i-1)}$ $m^{\text{f} i \text{fl}}$ 0 $m^{(i)}$ $\mathbf{A}_{2}^{(i)} \quad \mathbf{K}^{\mathfrak{t} \ i \, \mathfrak{t}} \quad \mathbf{A}_{1}^{(i)}$ 0 0 $\mathbf{A}_{2}^{(2)}$ $\mathbf{K}^{\mathfrak{t}} \, {}^{\mathfrak{c}\mathfrak{t}} \, {}^{\mathfrak{o}}\mathbf{A}_{1}^{(2)}$ $m^{\mathrm{f}\ 2\mathrm{f}}$ 0 $m^{(1)}$ $s(\omega)$

For layered media, we have a sparse global matrix

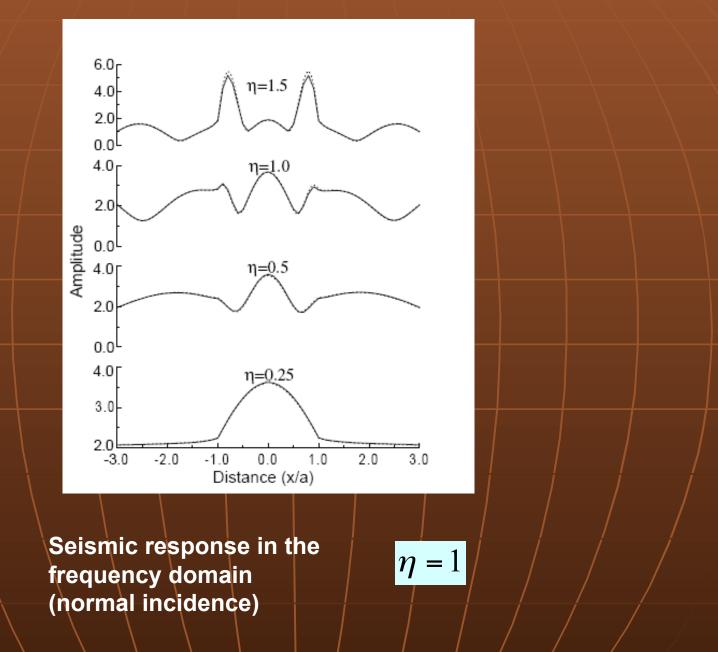


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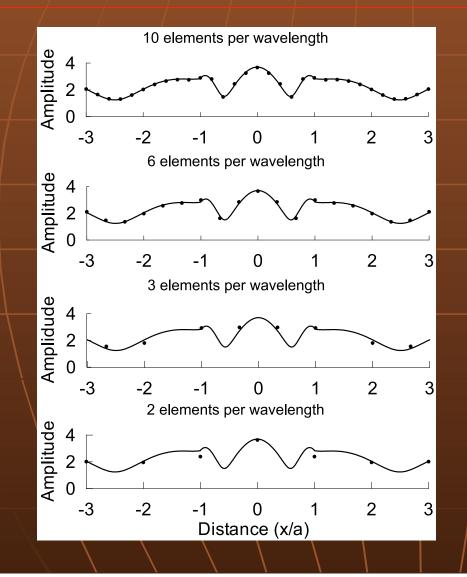


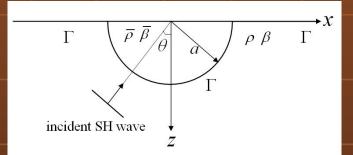
Benchmark test



Benchmark test

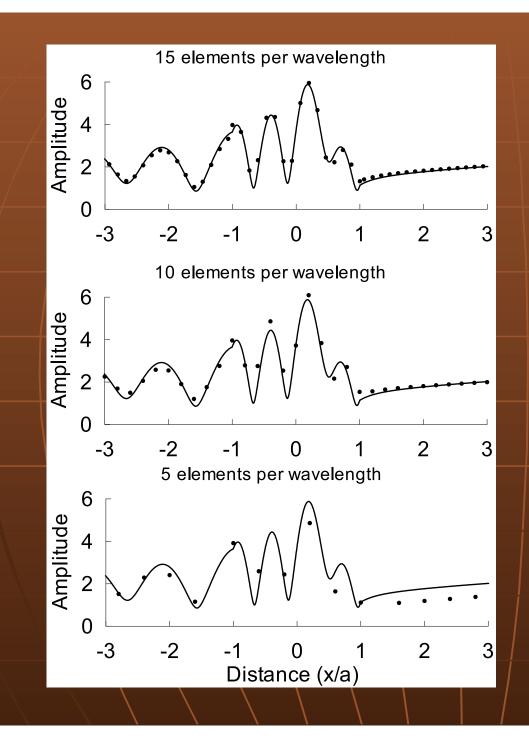
How many elements per wavelength do we need for seismic wave simulation?





Vertical incidence:

three elements per wavelength may be enough

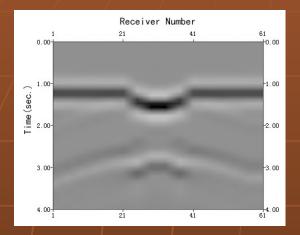


60° incidence:

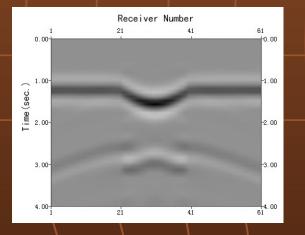
Five elements per wavelength still has some errors

If the geological structure is complex and we want to be more accurate, we should use more elements per wavelength

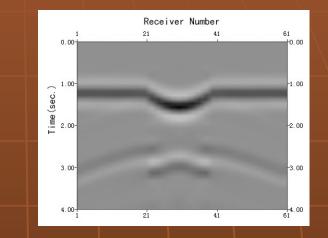
Benchmark test



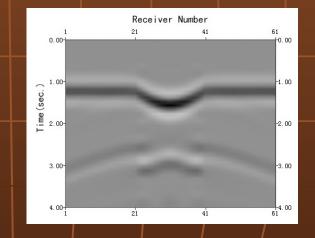
30 elements per wavelength



3 elements per wavelength



10 elements per wavelength



2 elements per wavelength

We see little difference from the synthetic data in the time domain

Benchmark test

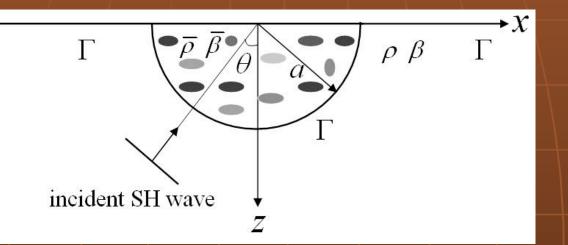
 In the frequency domain, we may have better chance to see the sensitivity of the parameters of the model.

•Seismic response in frequency domain is used to study ground motion in earthquake engineering.

•Before generating synthetic seismograms, we first do some analysis in the frequency domain.

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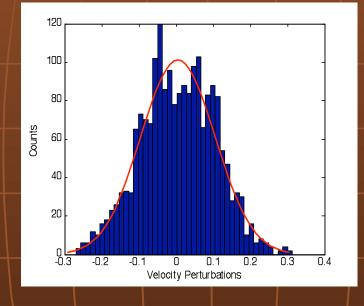
Modified valley model: random media in the valley

$$v(\mathbf{r}) = v_0(\mathbf{r}) + \delta v(\mathbf{r})$$

Six different random media

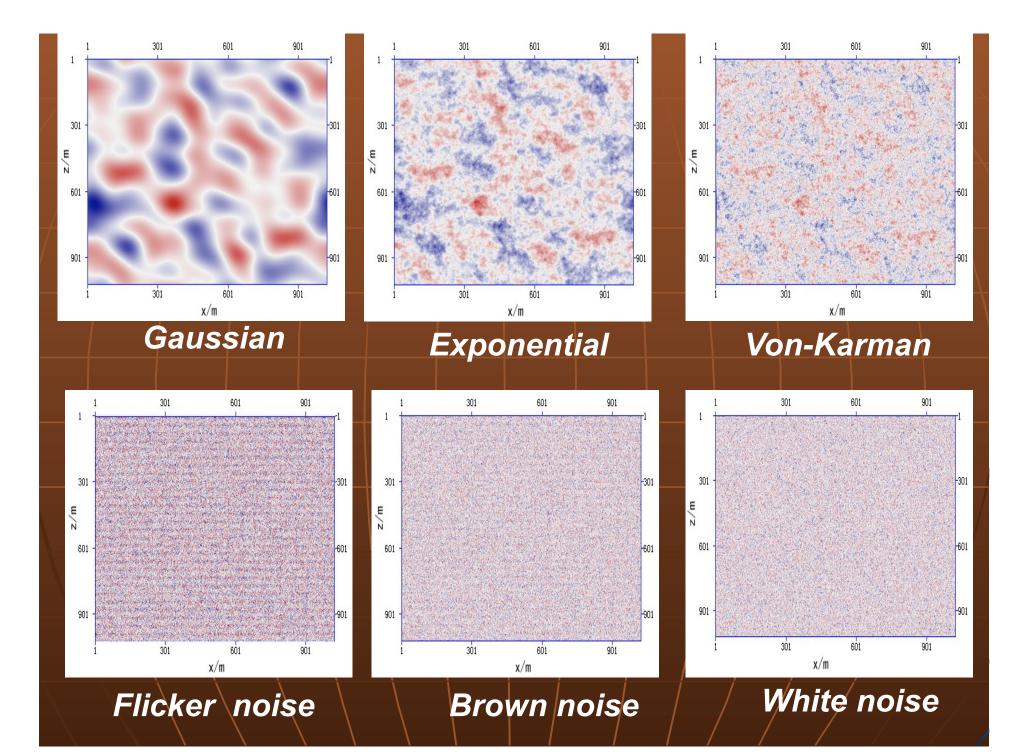
Different types of random media.	Filtering factor.	←
Gaussian	$\hat{f}_G(k) = \kappa \exp(-\frac{a_G^2 k^2}{8}) e^{-\frac{k^2}{8}}$	 a :autocorrelation length
Exponential.	$\hat{f}_{e}(k) = \kappa [a^{\bar{a}^{2}} + k^{2}]^{-\frac{d+1}{4}}$	N:Hurst index
Von-Karman 💩	$\hat{f}_{K}(k) = \kappa [a^{-2} + k^{2}]^{-\frac{d}{4} - \frac{N}{2}}$	
Flicker noise @	$\hat{f}_f(k) = \kappa k^{-1}$	
Brown noise₽	$\hat{f}_b(k) = \kappa k^{-3/2} \varphi$	
White noise.	$\hat{f}_{w}(k) = \kappa \checkmark$	Constant

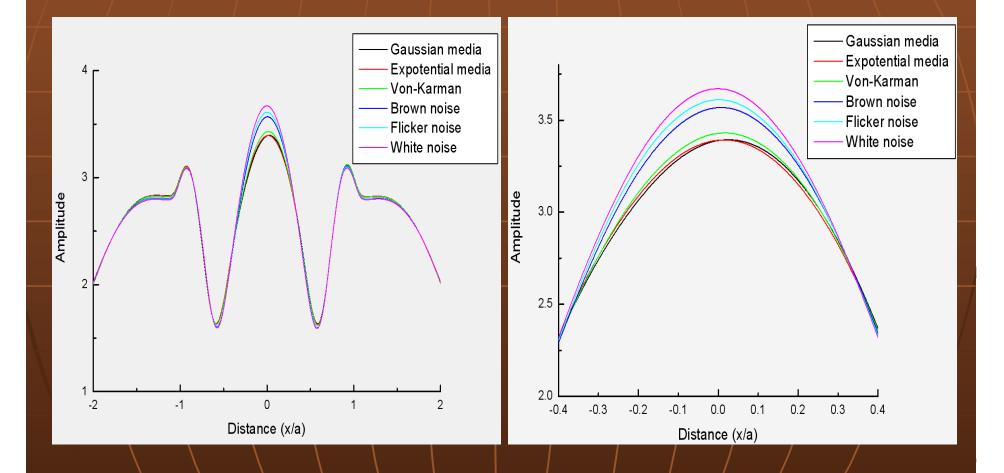
Gaussian type random media



a=2.5m, δ=0.1

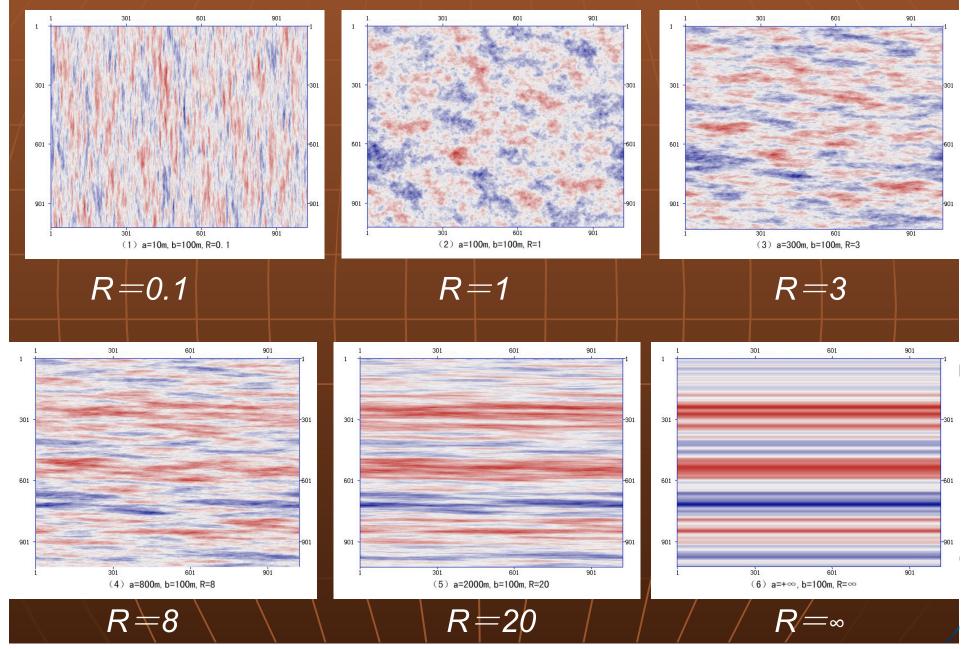
Although there are some deviations, the random media is generally assumed to satisfy statistical rules.



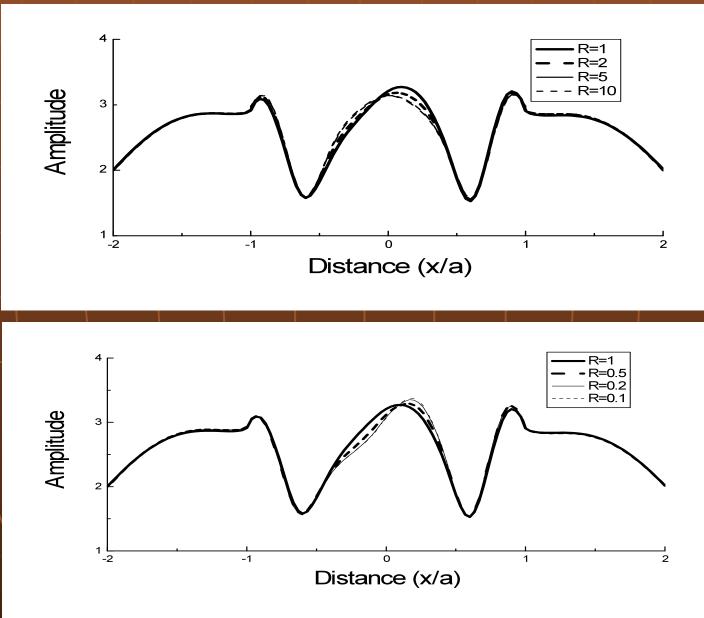


Frequency response with different types of random media

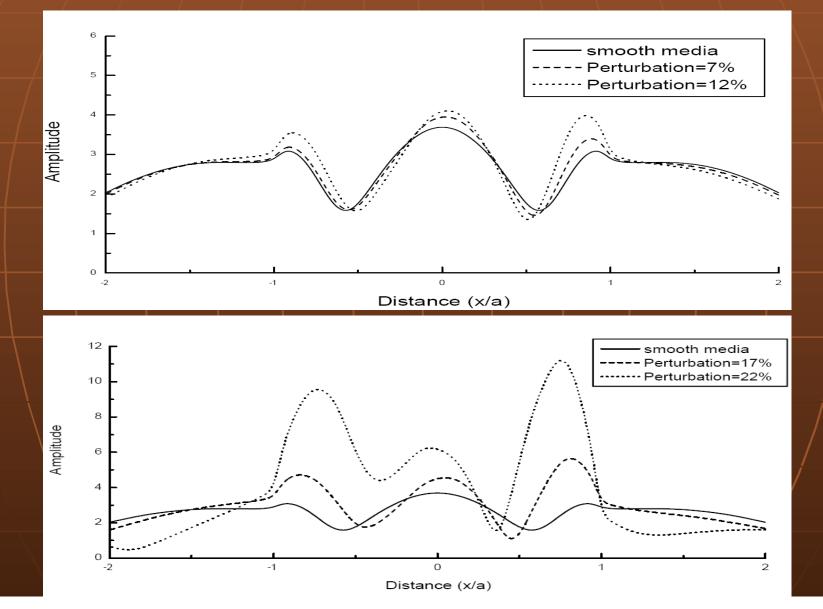
Different horizontal correlation length: R= horizontal correlation length/vertical correlation length



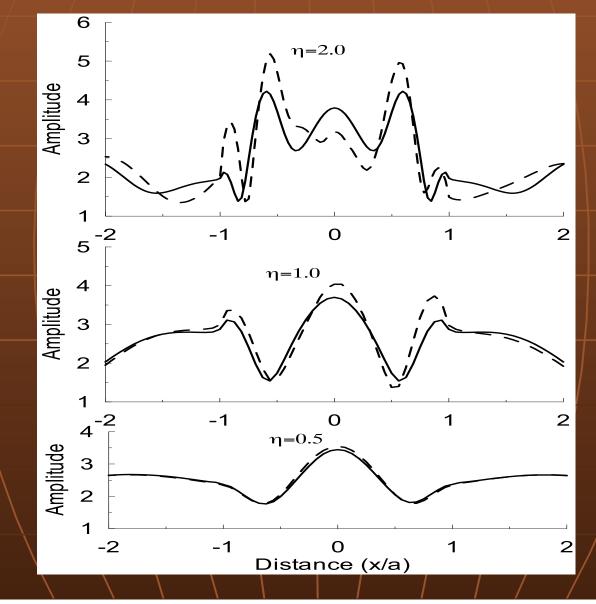
Wave scattering in the frequency domain Frequency response with different correlation lengths



Different Perturbations



Seismic response for different dimensionless frequencies



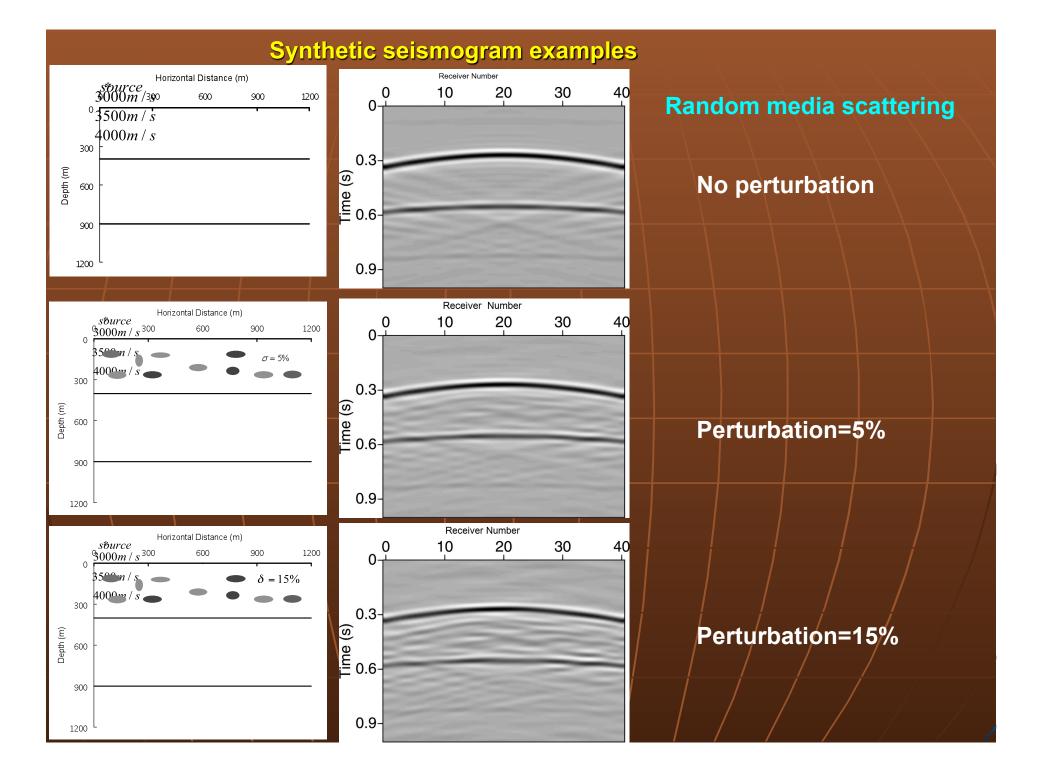
 Different types of random media and different parameters of the random media have different effect on the results.

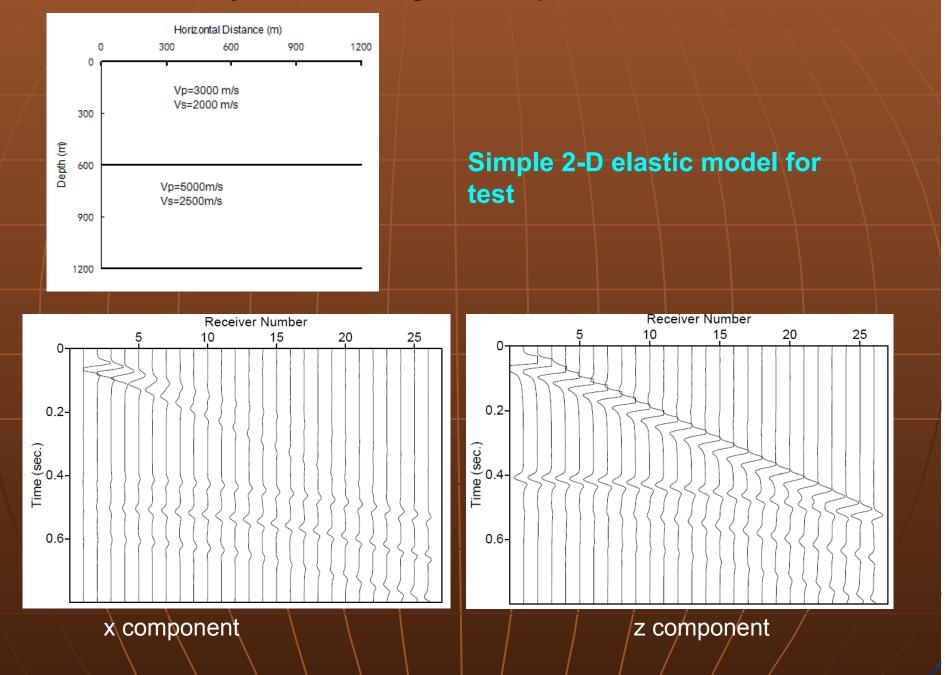
•Perturbations and dimensionless frequency (the complexity of the structure) affects the results much bigger than others.

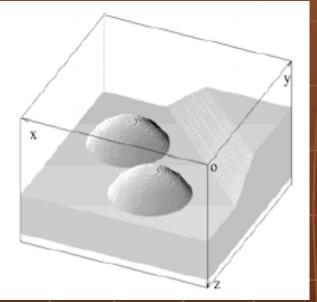
•The analysis allows us to better choose the parameters of the random media.

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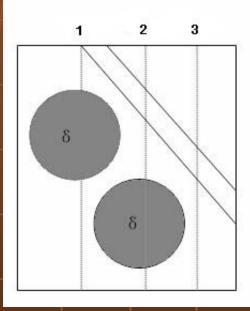
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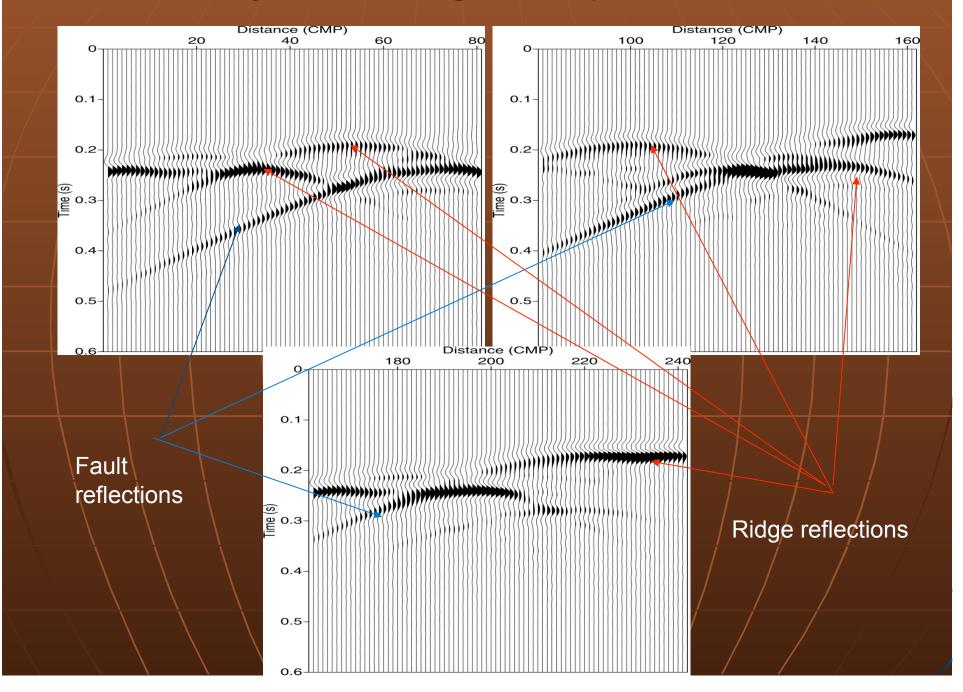
3D example: French Model

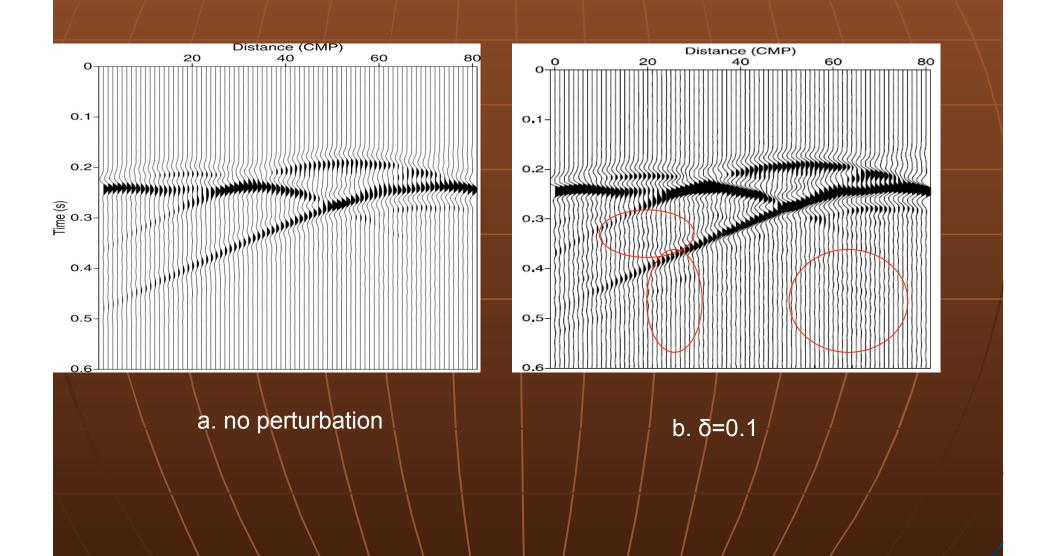


Model size: 625m*625m*100m

Radius of the hemisphere: 170.8m

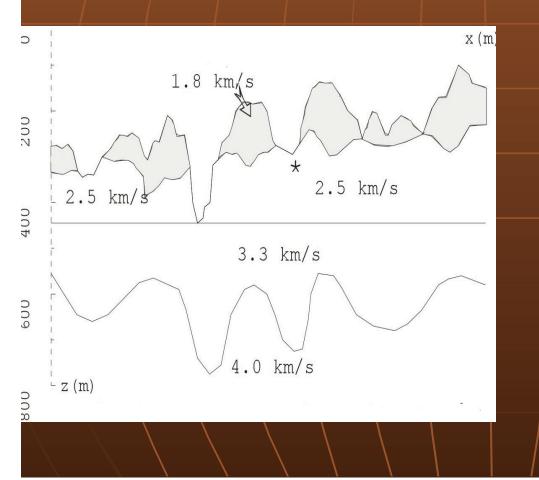
W.S. French 1974, Geophysics

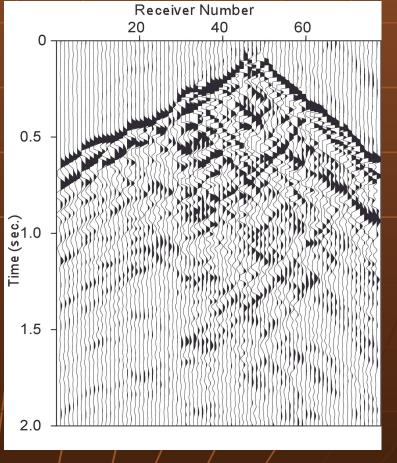




Irregular free surface to simulate land seismic acquisition.

Receivers are along the free surface





Conclusions

Characteristics of the modified Boundary Element Method:

- 1. Has the advantage of boundary element: requires less data to be calculated.
- 2. Explicitly use the boundary conditions of continuities across interfaces.
- 3. Easily adapt to curved free surface and irregular interfaces.
- 4. Seismic wave scattering by volume heterogeneities is included.5. While many practical problems need to be solved, this method provides an alternative for seismic forward modeling .

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